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New compactly supported wavelets for which both the scaling and wavelet functions have a high number of vanishing moments are presented. Such wavelets are a generalization of the so-called coiflets and they are useful in applications where interpolation and linear phase are of importance. The new approach is to parameterize coiflets by the first moment of the scaling function. By allowing noninteger values for this parameter, the interpolation and linear phase properties

suggested the design of orthonormal wavelet systems with vanishing moments for both the scaling and the wavelet functions. They were first constructed by Daubechies [9] and she named them *coiflets*.

In [1] shifted vanishing moments for the scaling function ϕ were used to obtain *one point quadratures*

$$\int_{\mathbb{R}} f(x) \phi(x - k) dx = \int_{\mathbb{R}} f(x) \phi'(x - k) dx; \quad (1.1)$$

where f is a sufficiently smooth function on the multiresolution space V_0 and $\int f(x) \phi(x - k) dx$ are good approximations of the coefficients of f in the expansion.

Since in [1] both matrices and operators were considered, the points $\{x_k\}$ were chosen to be $x_k \in kC$, where C is an integer. This “shift” corresponds to the first moment of the scaling function ϕ ,

$$\int_{\mathbb{R}} x \phi(x) dx; \quad (1.2)$$

Note that C is not the center of mass because $\int_{\mathbb{R}} x \phi(x) dx \neq 0$.

wavelet ψ . The key to our approach is to insist on a reasonable approximation to linear phase only in the passband of the associated low-pass filter m_0 .

It is well known that the properties defining coiflets can be easily described in terms of the coefficients $\{h_k\}$ of m_0 . The conditions on $\{h_k\}$ turn out to be dependent [14], and one of the goals of this article is to derive a system that is free of redundant equations. To obtain such a system, we perform a change of variables on $\{h_k\}$ via a linear transformation that has the shift τ as a parameter. This defining system is partly linear and partly quadratic. For filter lengths up to 20 the system can be explicitly solved via algebraic methods like

We also refer to such H as a QMF. As a consequence of (2.3), H satisfies the following functional equation:

$$H(z)/H(z^{-1}) = C H(z^{-1})/H(z)$$

In terms of the symbol H , (3.1) requires that

$$\sum_j 1^j j^k h_j \neq 0 \quad \text{for } 0 \leq k < M; \quad (3.3)$$

or equivalently, the factorization

$$H(z) \neq 0 \left(\frac{1+z}{2} \right)^M Q(z); \quad (3.4)$$

where $Q(z) \neq 0$.

As pointed out in the Introduction, we are interested in vanishing (shifted) moments of the scaling function

$$\mathcal{M}'_{;k} \neq 0$$

Equations (3.6) and (3.8) imply that the following four conditions, valid for all $k; 0 < k < N$, are equivalent:

$$\mathcal{M}'_k \int_{\mathbf{R}} x^{k'} \cdot x/dx \, D^{-k}; \tag{3.12}$$

$$\mathcal{M}'_{;k} \int_{\mathbf{R}} \cdot x \quad /^{k'} \cdot x/dx \, D^{-k_0}; \tag{3.13}$$

$$\mathcal{M}^h_{;k} \int_j \cdot j \quad /^k h_j \, D^{-k_0}; \tag{3.14}$$

$$\mathcal{M}^h_k \int_j j^k h_j \, D^{-k}; \tag{3.15}$$

Therefore, imposing moment conditions for either the wavelet or the scaling function amounts to finding a QMF H with moment conditions for its sequence of coefficients. In particular, the first moment of ψ , as defined in (1.2), equals the derivative of H at one,

$$\int \psi \, D H^0 \cdot 1/;$$

On the other hand, (2.1) forces $\int j H \cdot z/j \, D^{-1}$ for z on the unit circle. These last two properties allow us to show that the value $\int \psi$ should be within the support of ψ . Observe that this result is not evident since ψ is not a positive function.

PROPOSITION 3.2. *Let $H \cdot z/D$*

$$L^{-1} \cdot \int_{j \in \mathbb{D}^0} j^k h_j \, \mathbb{D}$$

Proof. Applying the operator $.xD^n$ (defined in the Appendix) at $z \square 1$ to the QMF equation (2.5), or taking derivatives at $\square 0$ in (2.2

its values on a shifted dyadic grid:

$$p_{x/D} = \sum_k P\left(\frac{Ck}{2^n}\right) \cdot 2^{-n} \quad k/:$$

Since at some scale any smooth function can be well approximated by polynomials, we have the almost interpolating property discussed in the Introduction.

Here we see the advantage of having both M

a neighborhood of $\mathbb{D} \setminus 0$

$$H.e^i \in \mathbb{D} a_H. \in e^{i\phi_H}. \in \quad \text{and} \quad 0. \in \mathbb{D} a_0$$

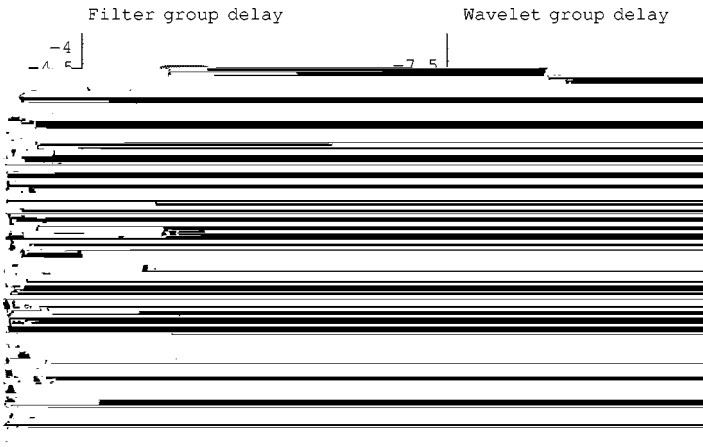


FIG. 2. Comparison between the group delays of m_0 and ψ . Maximal coefficient for wavelet: length 18, case b (top) and Daubechies' least asymmetric filter of the same length (bottom).

least asymmetric filter in Fig. 2 is defined as the maximally flat filter whose phase is as linear as possible within the whole band $\omega \in \omega_{\min}; \omega_{\max}$. See [8]

we obtain the *polyphase* equation

$$H_0(z)/H_0(z^{-1})/C H_1(z)/H_1(z^{-1})/D \frac{1}{2}. \tag{7.6}$$

The problem of finding a solution H of the QMF equation (2.5) is thus replaced by finding the solutions H_0 and H_1 of the *polyphase* equation. Instead of performing two operations on the variable z in (2.5), namely z and z^{-1} , in (7.6) we only have z^{-1} .

8. THE CONSTRUCTION OF COIFLETS

Recall that we can write any polynomial QMF as $H(z)/D \sum_{k=0}^L h_k z^k$, where $h_0, \dots, h_L \in \mathbb{C}$.

We describe a system for coiflets not in terms of $\{h_k\}$ but in terms of the new variables

$$a_k = \frac{1}{k!} \sum_j \binom{j}{2}^k h_{2j} \quad \text{and} \quad b_k = \frac{1}{k!} \sum_j \binom{j-1}{2}^k h_{2j+1};$$

where $0 \leq k \leq L$, and $D \frac{1}{2} L \leq 2L$. The transformation from $\{h_k\}$ to $\{a_k; b_k\}$ is linear and parameterized by \cdot . As before, $D \sum_j j h_j$ is the first moment of \cdot .

For what follows, it is more convenient to describe a_k and b_k for arbitrary $k \geq 0$, using the operator xD . We 2 k r i k r i k r l 0 n T r b i t i s (0 ,) - t (g 2 3 w 9 j / F 1 0 1 T

$$D \frac{1}{i^n} x D^n (x^{-2} H_0(x/x) \cdot x^{-2} H_0(x^{-1}/C x) \cdot x^{-1/2} H_1(x/x) \cdot x^{-1/2} H_1(x^{-1}/C x)) \cdot 1/$$

$$D \sum_{k=0}^n \frac{1}{k!} a_n \cdot a_k \cdot b_n \cdot b_k \tag{8.1}$$

for $0 \leq n \leq L - 2$.

If n is odd, the previous equation is always satisfied and then, as we remarked earlier, $\frac{L}{2}$ equations are enough to characterize a QMF of length L .

8.2. Linear Conditions

We now discuss how to rewrite the (linear) conditions (4.1) and (4.2) for coiflets in terms

8.3. A System for Coiflets

Combining (8.1) and (8.4), the system for coiflets can be written in terms of the unknowns $f; a_m; \dots; a_l; b_m; \dots; b_l$,

$$\frac{1}{2} \sum_{k=0}^{n-1} a_k \circ b_k \circ \dots \circ 1/k, a_n \circ k$$

8.3.3. *Coiflets with integer shifts.* Coiflets for *integer* choices of the shift were first computed by Daubechies [9]. In all cases that we computed, coiflets with integer shifts were always nonmaximal. In Table 1

An equivalent system, obtained via Gröbner bases where ϵ is treated as parameter is

function (by setting $a_3 \supset b_3$). In the latter case, (9.3) becomes

f

TABLE 2
Coiflet Filters of Length 8

	k	h_k		k	h_k
<i>MD</i> 3	0	0.00899863735774892	<i>MD</i> 3	0	0.03952785122359428
<i>ND</i> 5	1	0.02054552466216258	<i>ND</i> 5	1	0.1271031281675352
<i>D</i> 1	2	0.2202099211463259	<i>D</i> 2	2	0.5323389066059403
Case Na	3	0.5701914465849665	Case Nb	3	0.440002251136967
MAXIMAL	4	0.3422577968313942	MAXIMAL	4	0.005981694132267174
	5	0.07306459213264614		5	0.07120132136770919
	6	0.05346908061997128		6	0.01317063874992116
	7	0.02341867020984207		7	0.004095942063206933
<i>MD</i> 3	0	0.1646660519380485	<i>MD</i> 3	0	0.3040839480619514
<i>ND</i> 3	1	0.5074101320413008	<i>ND</i> 3	1	0.414464867958699
<i>D</i> 1	2	0.4435018441858542	<i>D</i> 1	2	0.02524815581414562
Case 1a	3	0.02223039612390291	Case 1b	3	0.2566053961239029
	4	0.1310018441858543	BAD	4	0.2872518441858542
	5	0.02223039612390291		5	0.2566053961239029
	6	0.02283394806195145		6	0.1165839480619514
	7	0.007410132041300974		7	0.085535132041301
<i>MD</i> 3	0	0.01938529090153145	<i>MD</i> 3	0	0.0850102909015314
<i>ND</i> 3	1	0.1854738954507657	<i>ND</i> 3	1	0.1332761045492342
<i>D</i> 2	2	0.5581558727045942	<i>D</i> 2	2	0.2449691272954056
Case 2a	3	0.3810783136477028	Case 2b	3	0.5376716863522972
	4	0.05815587270459436	UGLY	4	0.2550308727045943
	5	0.06857831364770281		5	0.2251716863522971
	6	0.01938529090153145		6	0.0850102909015314
	7	0.002026104549234272		7	0.05422389545076572
<i>MD</i> 3	0	0.05191993211769211	<i>MD</i> 3	0	0.01058006788230788
<i>ND</i> 3	1	0.0234375	<i>ND</i> 3	1	0.0234375
<i>D</i> 3	2	0.3432597963530763	<i>D</i> 3	2	0.2192402036469236
Case 3a	3	0.5703125	Case 3b	3	0.5703125
	4	0.2192402036469236		4	0.3432597963530763
	5	0.0703125		5	0.0703125
	6	0.01058006788230788		6	0.05191993211769211
	7	0.0234375		7	0.0234375

Note. $_1D$ 2.977273091796802, $_2D$ 2.239549738364678.

9.2. Coiflets of Length 18

A similar analysis can be done for filters of length 18. In Table 4, we present a summary of our findings by listing the filter coefficients for two cases: coiflets with integers shifts and maximal coiflets. Filter coefficients are listed in Table 6.

Even at higher numbers of vanishing moments and different lengths, we still found UGLY and BAD filters. They always correspond to coiflets with integer shifts, but it is not a peculiarity of that case. Varying α , we found regions of nonmaximal coiflets with a similar behavior.

TABLE 3
Summary of All Maximal Coiflets and Coiflets with Integer Shifts for Length 8

Filter	M	N	$k_Q k$	2^{M-1}	Remarks		
Na	2.97727	3	5	1.45584	2.8764	4	
Nb	2.23955	3	5	1.44599	2.94511	4	
Ma	1.00539	4	3	1.77557	5.91608	8	Daubechies' Extremal Phase
Mb	2.98547	4	3	1.77557	5.91608	8	Daubechies' Least Asymmetric
1a	1	3	3	1.77528	2.16403	4	
1b	1	3	3	0.14666	14.9356	4	BAD
2a	2	3	3	1.42232	3.11099	4	
2b	2	3	3	0.93596	6.91099	4	UGLY
3a	3	3	3	1.77341	2.16473	4	
3b	3	3	3	1.46353	2.82288	4	

Note. Coefficients are listed in Table 2. The maximal case for wavelets coincides with Daubechies' maximally flat filters.

In Figs. 11 and 12, we plotted $|m_0|$ and ρ for the cases 6c (UGLY) and 5b (BAD) with length 18. The cases 7d and 6d, as listed in Table 4, exhibit a similar behavior. Even though their filter moduli do not oscillate as much as their counterparts of length 8, their behavior is clearly different than those for which $k_Q k$ remains below 2^{M-1} . As an example of the latter situation, consider the filter 7c. The associated wavelet has only six vanishing moments, but its Sobolev exponent is higher than the exponent for Daubechies' wavelets which have nine vanishing moments.

Note that in all the plots for wavelets in the Fourier domain, the support of the functions is actually wider than shown.

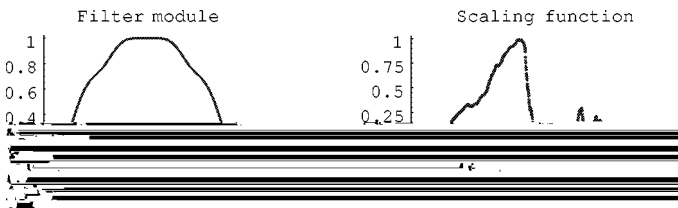


FIG. 3. Integer shift coiflet: length 8, shift 2, case b (UGLY). Plots of absolute value of filter m_0 and scaling function.



FIG. 4. Integer shift coiflet: length 8, shift 1, case b (BAD). Plots of absolute value of filter m_0 and scaling function.

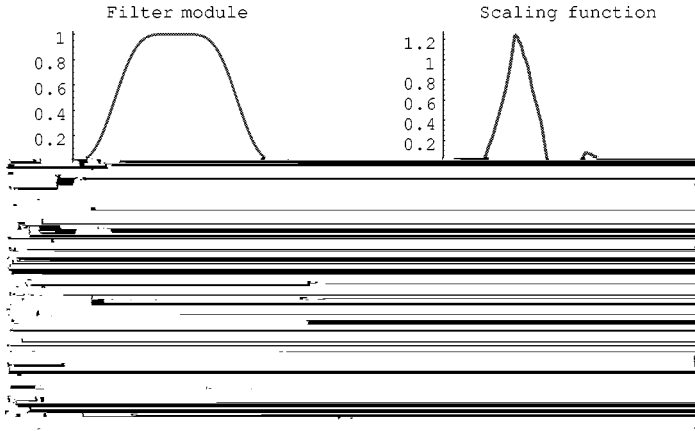


FIG. 5. Integer shift coiflet: length 8, shift 2, case a. Plots of scaling function and filter m_0 .

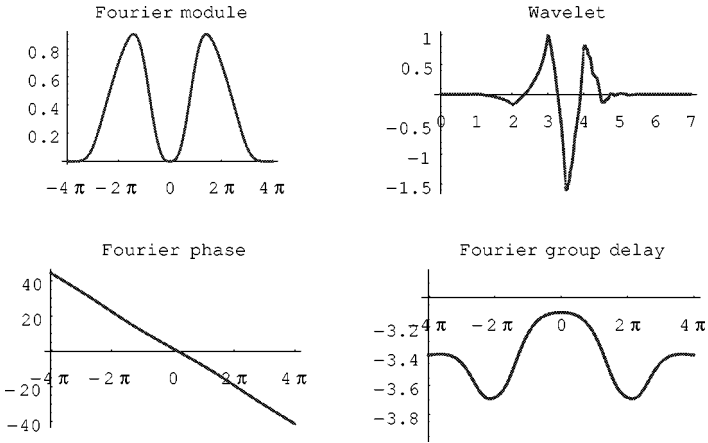


FIG. 6. Integer shift coiflet: length 8, shift 2, case a. Plots of wavelet function in both time and Fourier domain (absolute value, phase, and group delay).

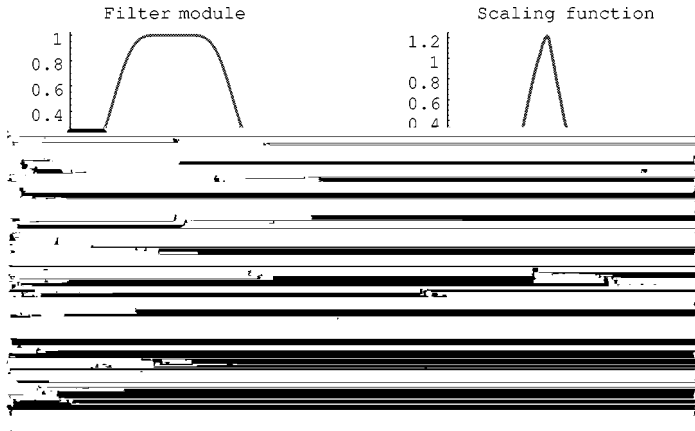


FIG. 7. Integer shift coiflet: length 8, shift 3, case a. Plots of scaling function and filter m_0 .

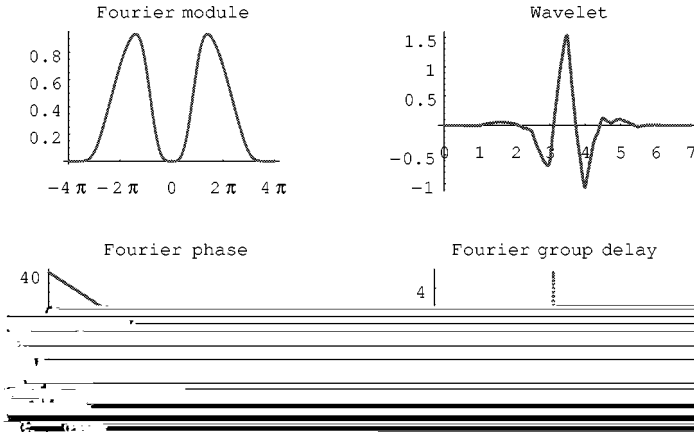


FIG. 8. Integer shift coiflet: length 8, shift 3, case a. Plots of wavelet function in both time and Fourier domain (absolute value, phase, and group delay).

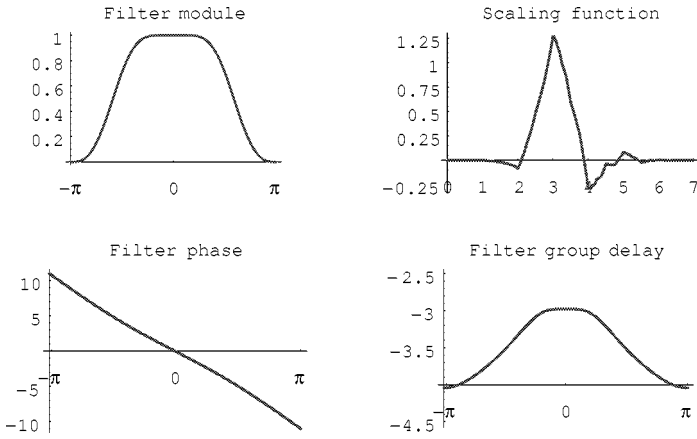


FIG. 9. Maximal coiflet for scaling function: length 8, shift 2.9773. Plots of the scaling function and filter m_0

TABLE 4
Summary of All Maximal Coiflets, Coiflets with Integer Shifts, and Two Daubechies' Maximally Flat Filters for Length 18

Filter		M	N		$\kappa_Q \kappa$	2^{M-1}	Remarks
Na	7.81041	6	9	2.5149	16.5942	32	Listed in Table 5
Nb	7.1771	6	9	2.49853	17.2438	32	Listed in Table 5
Ma	5.94301	7	7	2.74543	33.9874	64	Listed in Table 5
Mb	4.5681	7	7	2.71944	36.2534	64	Listed in Table 5

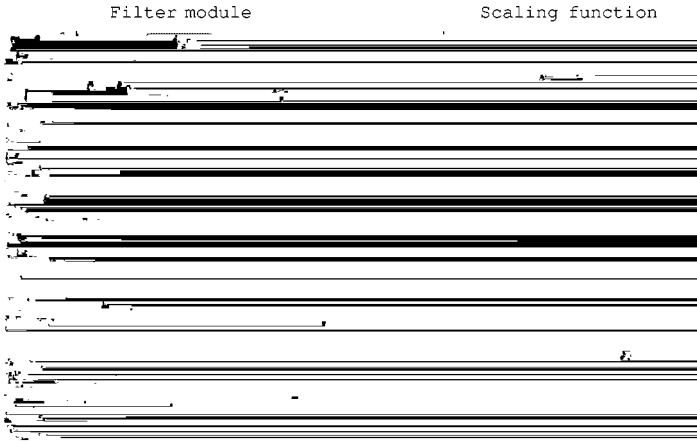


FIG. 13. Maximal coiflet for scaling function: length 18, shift 7.1771. Plots of filter m_0 and scaling function.

10. CONCLUSION

The approach taken in this paper allows one to construct and classify coiflets, which are wavelets with a high number of vanishing moments for both the scaling and wavelet functions. Coiflet filters are useful in applications where interpolation and linear phase are of particular importance.

We introduced a new system for FIR coiflets. In all cases investigated, the system had a minimal set of defining equations. For filters of length up to 20, the system can be solved explicitly, and the filter coefficients can thus be accurately determined. For longer filters we applied numerical methods to compute some solutions. For a few specific examples we studied the properties of coiflets corresponding to both integer and noninteger values of the first moment of the scaling function. Nevertheless, the problem of the existence of coiflet filters of arbitrary length and their full classification remains open.

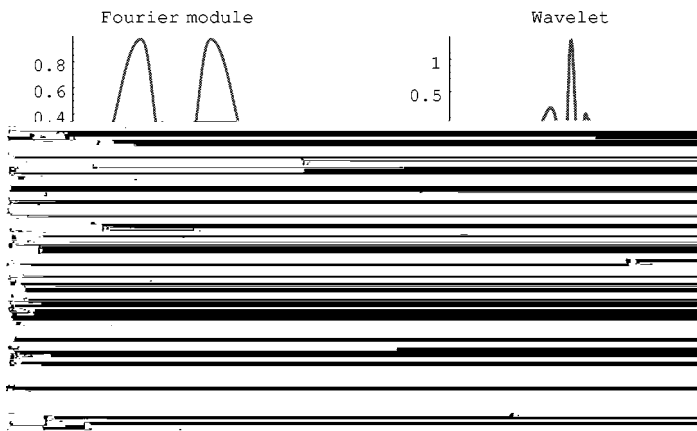


FIG. 14. Maximal coiflet for scaling function: length 18, shift 7.1771. Plots of wavelet function in both time and Fourier domain (absolute value, phase, and group delay).

TABLE 5
Coiflet Filters of Length 18: Maximal Case

	k	h_k		k	h_k
MD_6	0	0:00006423105557385401	MD_6	0	0:0002036914946771235
ND_9	1	0:0002979447888413989	ND_9	1	0:0002488151932121008
D_1	2	0:0004927238418624587	D_2	2	0:00221156402899935
Case Na	3	0:004159721116204626	Case Nb	3	0:005347581803838808
	4	0:001356751057023208		4	0:02049652597342785
	5	0:03424128516618039		5	0:034353284830D 47264271 00

From (A.1) and (A.3)

$$z^n D^n f(z) / \mathbb{D} = \sum_{k=0}^n s_k^n \cdot x D^k f(z) / \mathbb{D} \quad (\text{A.4})$$

Note that for a polynomial of degree r , it is not true that $x D^n P(z) / \mathbb{D}$ is zero for $n > r$. However, these values are linear combinations of $x D^k P(z) / \mathbb{D}$ for $n \leq r$, as we show in the

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